

CONTENT

I. Geophysical context

II. Inverse modelling

III. A new finite element solver

IV. Virtual inverse rheometry

- ~~Reliability of surface velocity observations~~
- ~~Power-law index identification~~
- Fluid consistency: sensitivity
- Fluid consistency: identification

V. The basal slipperiness

VI. Conclusions & Perspectives

VIRTUAL INVERSE RHEOMETRY

- Power-law description : $\eta(\mathbf{u}) = \eta_0(T) \|\underline{\mathbf{D}}(\mathbf{u})\|_F^{s-2}$
→ uncertain parameters:
Exponent s
Consistency η_0
Temperature $T(z)$: linear decrease from bottom to surface
- « Ill-posed »: $\eta(\mathbf{u}) \longrightarrow +\infty$ when $\|\underline{\mathbf{D}}(\mathbf{u})\| \rightarrow 0$ since $s - 2 < 0$
→ singularity at free surface (vanishing shear)

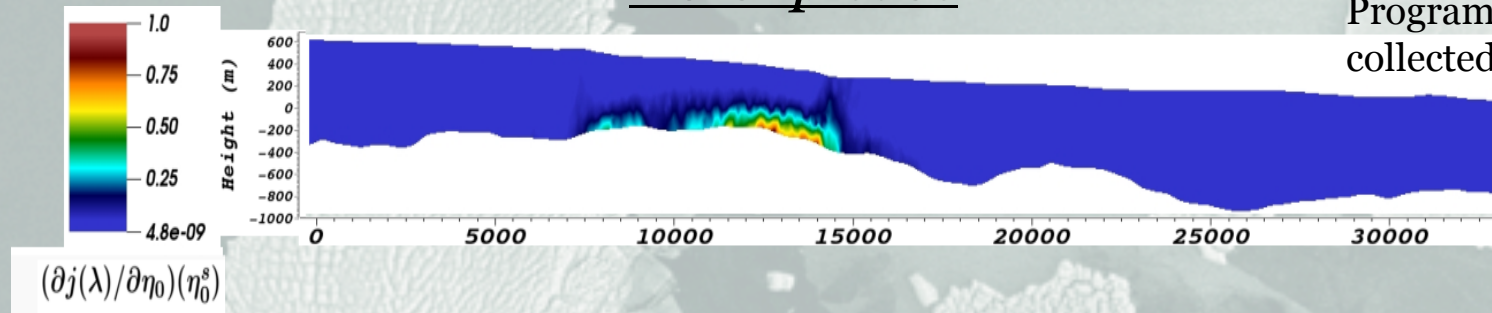
Goal:

- Determine the importance of rheological parameters
- Provide calibration/identification

Fluid consistency $\eta_0(T)$: sensitivity

No-slip case

Program ICECAP, 2010,
collected by B. Legrésy



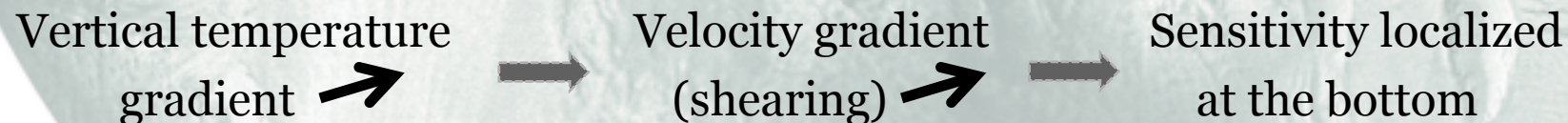
- High sensitivity = highly sheared area
 - ➡ localized at the accelerating zone (bump)
 - ➡ concentrated at bottom (no-slip)

Behaviour at high shear-rate dominant

- Sensitivity independant of the vertical temperature gradient

Sliding case

- Relaxation of the high shearing at bottom ➡ distribution of the shear-rate

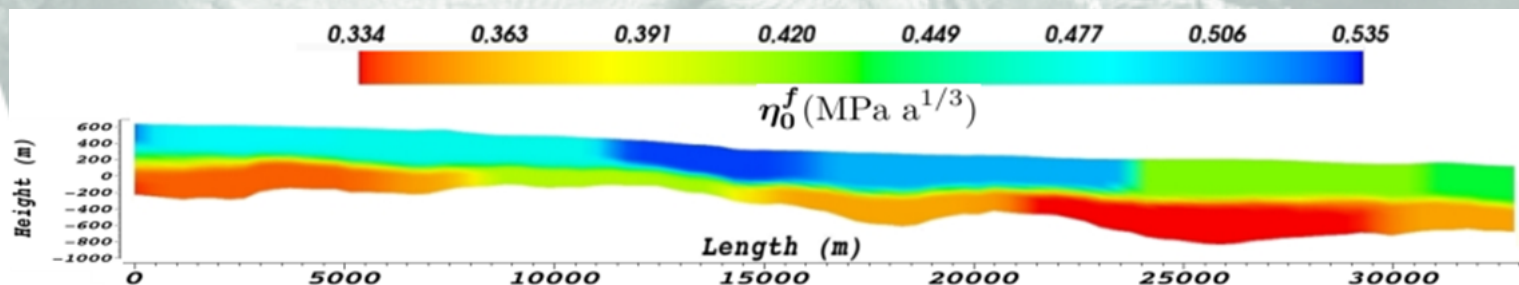


Fluid consistency $\eta_0(T)$: identification

- $\eta_0(T)$: spatially distributed value (to be inferred)
➔ requires regularization (Tykhonov)
- Absolute sensitivity: low compared to the exponent one
- Cost function :

$$j(\eta_0) = \int_{\Gamma_s} \|u_s^{obs} - u_s(\eta_0)\|_2^2 d\mathbf{x} + \gamma_1 \int_{\Gamma_s} \|\partial_x \eta_0\|_2^2 ds + \gamma_2 \int_{\Gamma_s} \|\partial_z \eta_0\|_2^2 ds$$

- Regularization: *good a priori* required on the temperature gradient
Given exact ratio: $\gamma_1 = \gamma_2/7$, Initial-guess : surface value



Synthetic surface
data: 0.1% noise
Final error ~ 9%

Otherwise ill-posed inverse problem, converge to wrong optimal value

Results presented in Martin-Monnier, J. Geophysical Research, submitted

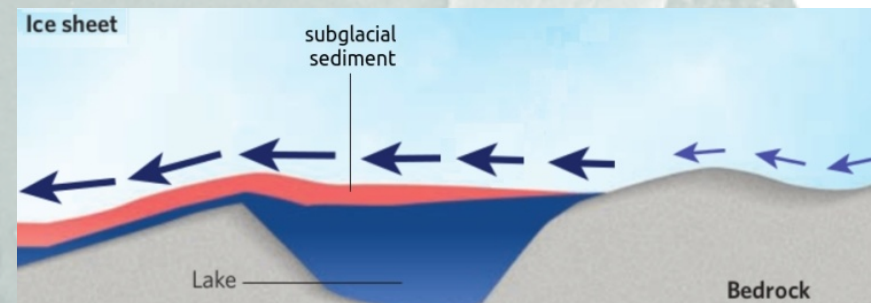
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- I. Geophysical context
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- V. The basal slipperiness**
 - Sensitivity to a free-slip area
 - Sensitivity and slip ratio
 - Formulation of an incomplete adjoint
 - Friction coefficient invertibility
- VI. Conclusions & Perspectives

THE BASAL SLIPPERINESS

- Major component of ice dynamics
 - *models heterogeneous interfacial layer*

Regelation, enhanced creep, subglacial hydrology, basal sediment deformation



$$|\tau_b|^{m-1} \tau_b = \beta u_b^\tau$$

1 parameter involving various processes

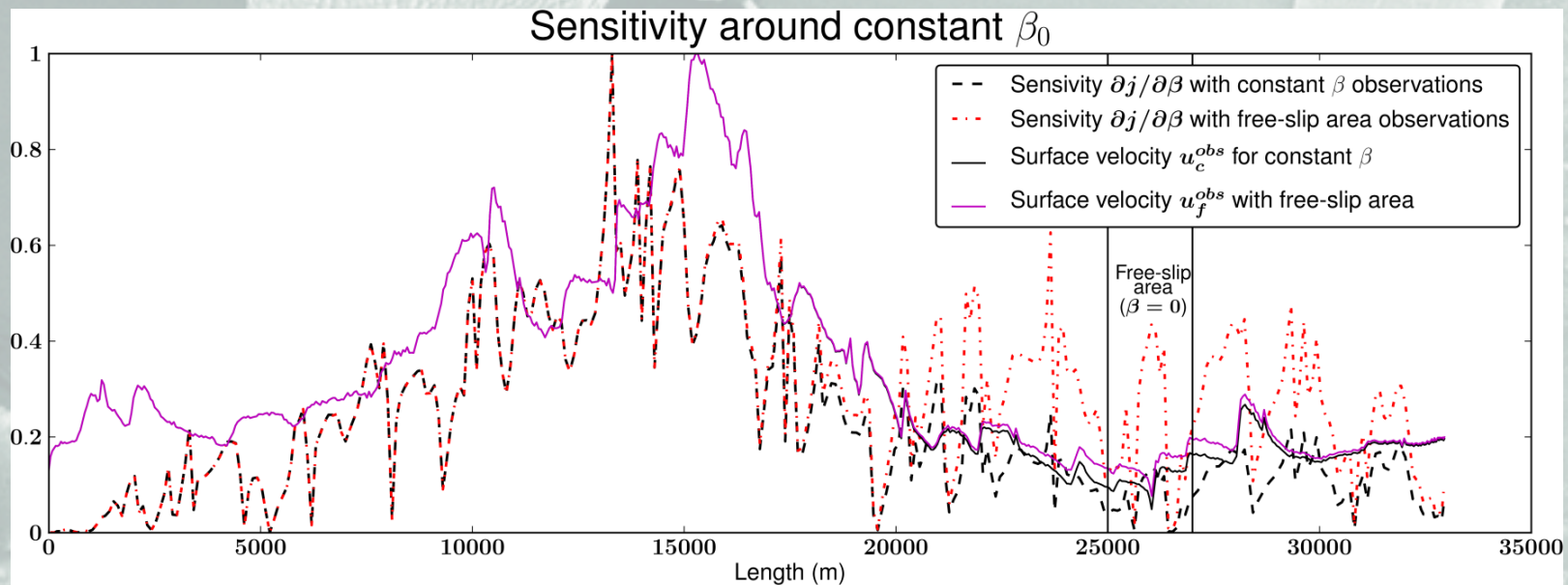
- Generally unmeasurable

Goal:

- ➔ Sensitivity analysis: understand β modelling role
- ➔ Identification: infer the *bedrock-ice interaction* properties

Sensitivity to a free-slip area

Experiment 1: effect of a local free-slip on the sensitivity to friction



- Sensitivity correlated to the surface velocities (hence the topography)
- Local free-slip area (2km) $\longrightarrow$ large high sensitivity area ($\sim 10\text{km}$)
 $\longrightarrow$ almost no effect on u_s
- Transmission of basal movement to the surface: **filtered** and **non-local**

Sensitivity to a free-slip area

- *Quasi-static sensitivities:*
given a snapshot (surface & bedrock topography + surface velocities)
→ help to see beyond filtering and non-local
- *Free-slip area:*
possibly local in space and time (e.g. fluctuation of subglacial meltwater pressure)

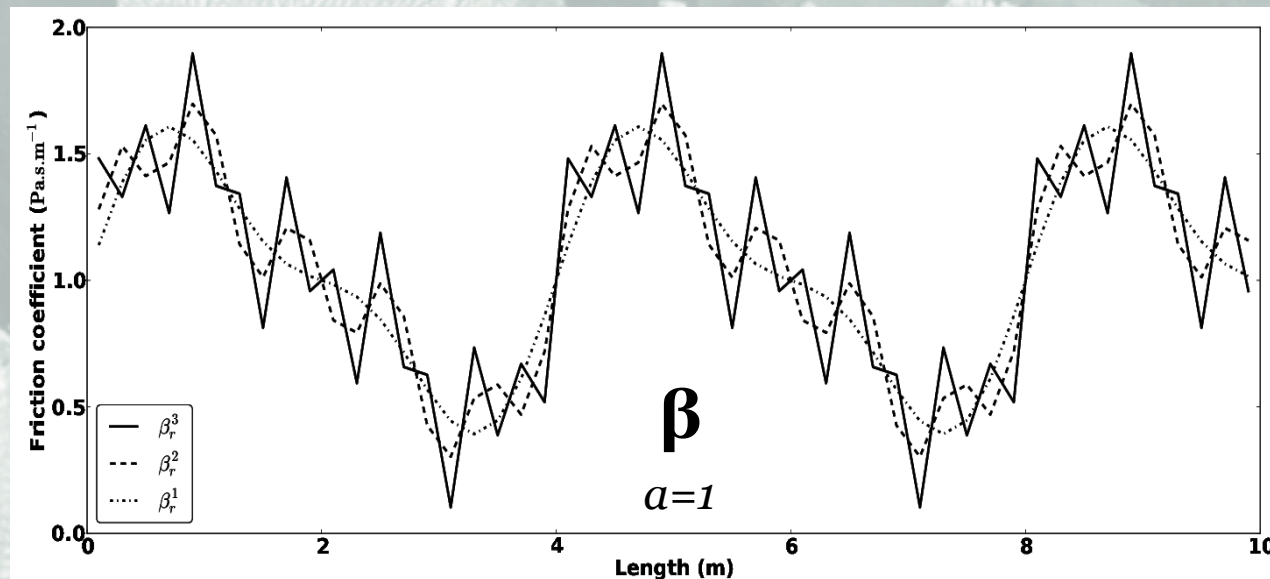


Sensitivity analysis could allow to *detect momentary increase* in sliding rate (e.g. « *stick-slip* » phenomena)

Sensitivity and slip ratio r

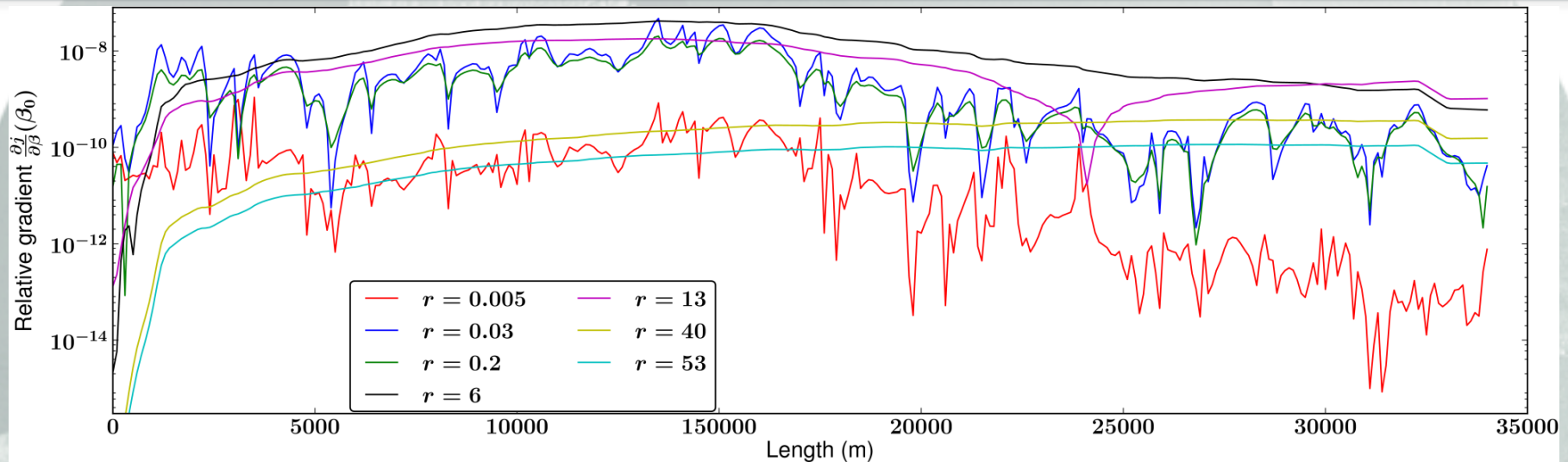
Experiment 2: effect of the slip ratio on the sensitivity to friction

- Slip ratio $r = \overline{u_b} / |\overline{u_s} - \overline{u_b}|$: balance between deformation and sliding
- Various frequencies in the friction coefficient β



$$\beta_r^n(x) = a + \frac{a}{2} \sin\left(\frac{2\pi x}{\lambda}\right) + \frac{a}{5} \sum_{i=1}^n f_i(x) \quad \beta = \beta_r^3$$

Sensitivity and slip ratio r



Sensitivity to the friction coefficient β for various slip ratio r

- *Low slip ratio: $r < 0.01$:*
high frequencies captured
low absolute values
➔ filtering effects
- *High slip ratio: $r > 5$*
low frequency information
average behaviour only
➔ non-local effects

Best situations to perform identification: *intermediate sliding*

Formulation of an incomplete adjoint

Recall: *adjoint model*

algorithmic differentiation of a fixed point type *iterative routine*
(Stokes power-law): $y = \Phi(y, u)$

Reverse Accumulation (Griewank et al., 1993)

1) derivation of the *computational graph*
(node by node, one node n_i per function value Φ)

→ one adjoint quantity per node = $\frac{\partial \Phi}{\partial n_i}$

2) compute the adjoint state in reverse order

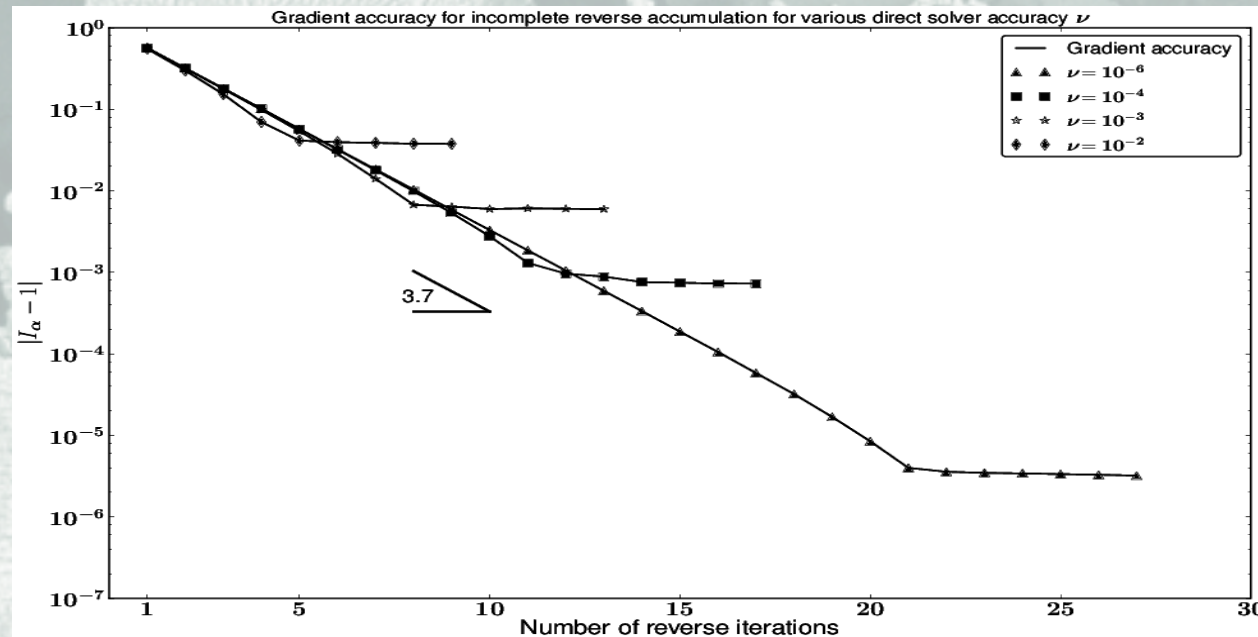
→ adjoint state = $\sum_i \frac{\partial \Phi}{\partial n_i}$ (chain formula)

Corollary: Adjoint state precision = Direct state precision
Christianson, 1992 **(~fixed point threshold)**

Formulation of an incomplete adjoint

Idea: retain « some » of the $\sum_i \frac{\partial \Phi}{\partial n_i}$ to compute the adjoint state
from 1 (« self-adjoint ») to all (complete adjoint)

What is the precision of the resulting *incomplete adjoint*?



➔ Adjoint precision adjustable a priori from data accuracy & direct solver precision

Friction coefficient invertibility

Experiment 3: Assess the limits of identifiability of the friction coefficient in Fourier space

- level of captured frequencies (in terms of thickness h)
- reconstructed amplitude

Setup:

- Synthetic data: 1% random noise
- Density: 1km (~ 1 ice thickness h) between each data point
- Realistic topography
- Non-linear friction law: $|\tau_b|^{m-1} \tau_b = \beta u_b^\tau$ with $m=3$

➔ *Comparison of exact adjoint and « self-adjoint » method*

Friction coefficient invertibility

Slip ratio: $r=0.005$ (strong friction)

- *Complete adjoint:*

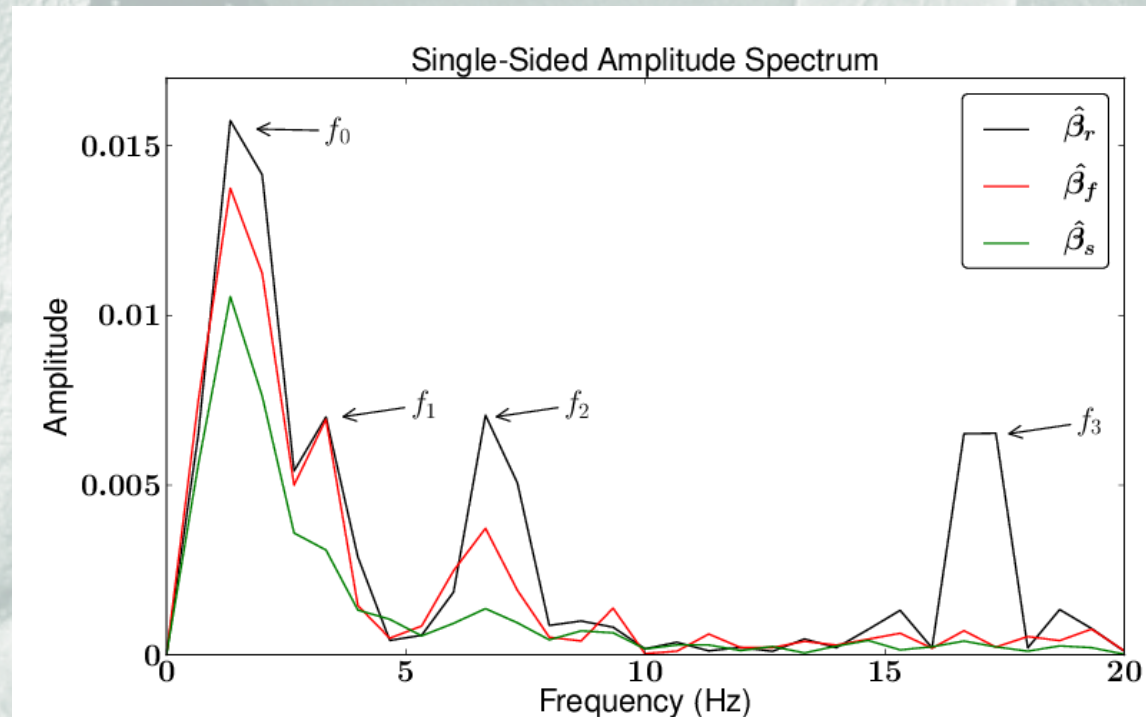
High frequencies captured
Smaller amplitude

Filtering dominant

- « *Self-adjoint* »:

carrier frequency only

➡ not robust at
low slip ratio



	f_0	f_1	f_2	f_3
Wavelengths w.r.t $h = 1\text{km}$	$20h$	$10h$	$5h$	$2h$
Wavelengths w.r.t $dx = 100\text{m}$	$200dx$	$100dx$	$50dx$	$20dx$
Wavelengths w.r.t $L = 33.3\text{km}$	1.66Hz	3.33Hz	6.66Hz	16.6Hz

Friction coefficient invertibility

Slip ratio: $r=0.5$ (intermediate sliding)

- *Complete adjoint:*

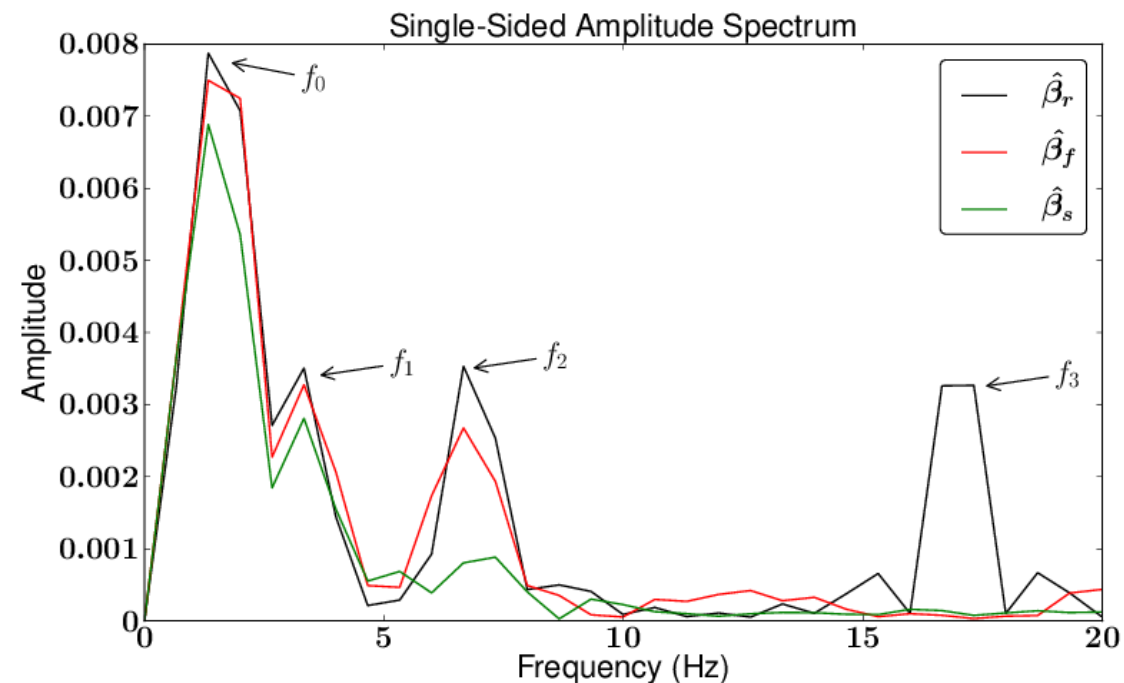
accurate up to
5h wavelengths

- « *Self-adjoint* »:

accurate up to
10h wavelengths

Complete adjoint

more accurate



	f_0	f_1	f_2	f_3
Wavelengths w.r.t $h = 1\text{km}$	$20h$	$10h$	$5h$	$2h$
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Wavelengths w.r.t $L = 33.3\text{km}$	1.66Hz	3.33Hz	6.66Hz	16.6Hz

Friction coefficient invertibility

Slip ratio: $r=50$ (rapid sliding)

- *Complete adjoint:*

Amplitude well represented
High frequencies not captured

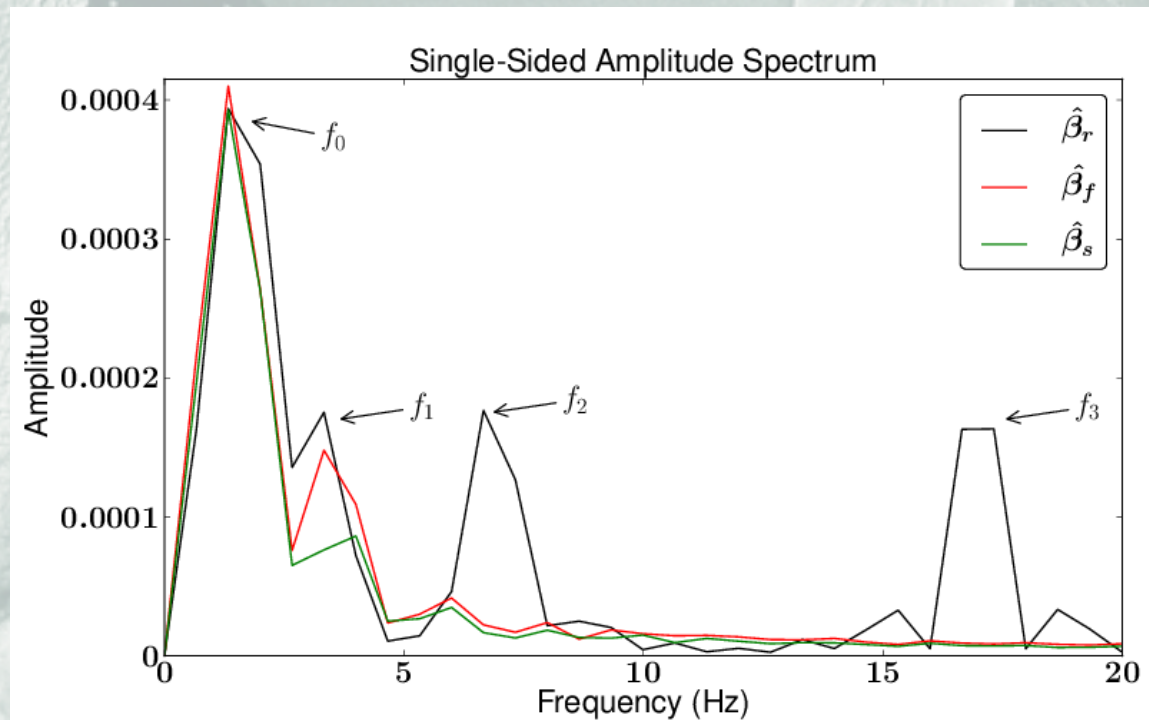
Non-local effects dominant

accurate up to 10h
wavelengths

- « *Self-adjoint* »:

Only the carrier wave

accurate up to 20h
wavelengths



	f_0	f_1	f_2	f_3
Wavelengths w.r.t $h = 1\text{km}$	$20h$	$10h$	$5h$	$2h$
Wavelengths w.r.t $dx = 100\text{m}$	$200dx$	$100dx$	$50dx$	$20dx$
Wavelengths w.r.t $L = 33.3\text{km}$	1.66Hz	3.33Hz	6.66Hz	16.6Hz

Results presented in *Martin-Monnier, The Cryosphere, submitted*

Overview

3) Basal sliding

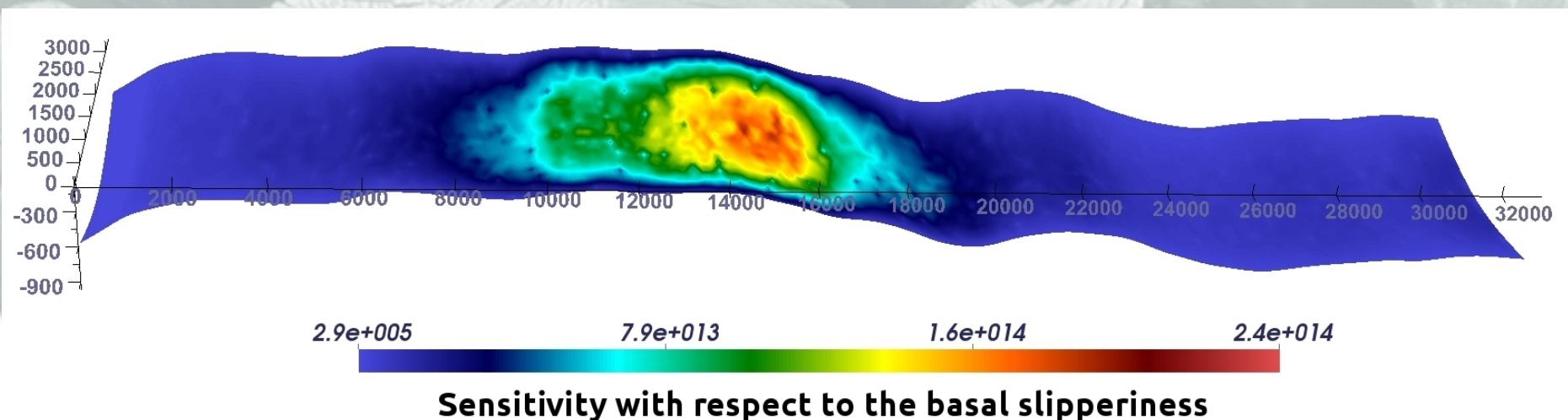
- high frequency events captured ($\sim 2x$ thickness)
- detection of momentary basal slipperiness variation (« stick-slip »)
Martin-Monnier, SIAM-SISC, submitted
- Adjustable incomplete adjoint (from data accuracy & direct solver precision)
- Assessment of transmission of basal variability:

<i>Friction</i>	Strong	Intermediate	Low
<i>Slip ratio</i>	Small	$\sim$ unity	Higher
<i>Sensitivity</i>	low	high	smooth
<i>Transmission of basal variability</i>	filtered $\sim 5h$ (approx. amplitude)	« optimal » $\sim 5h$	non-local $\sim 10h$

Martin-Monnier, The Cryosphere, submitted

Some perspectives

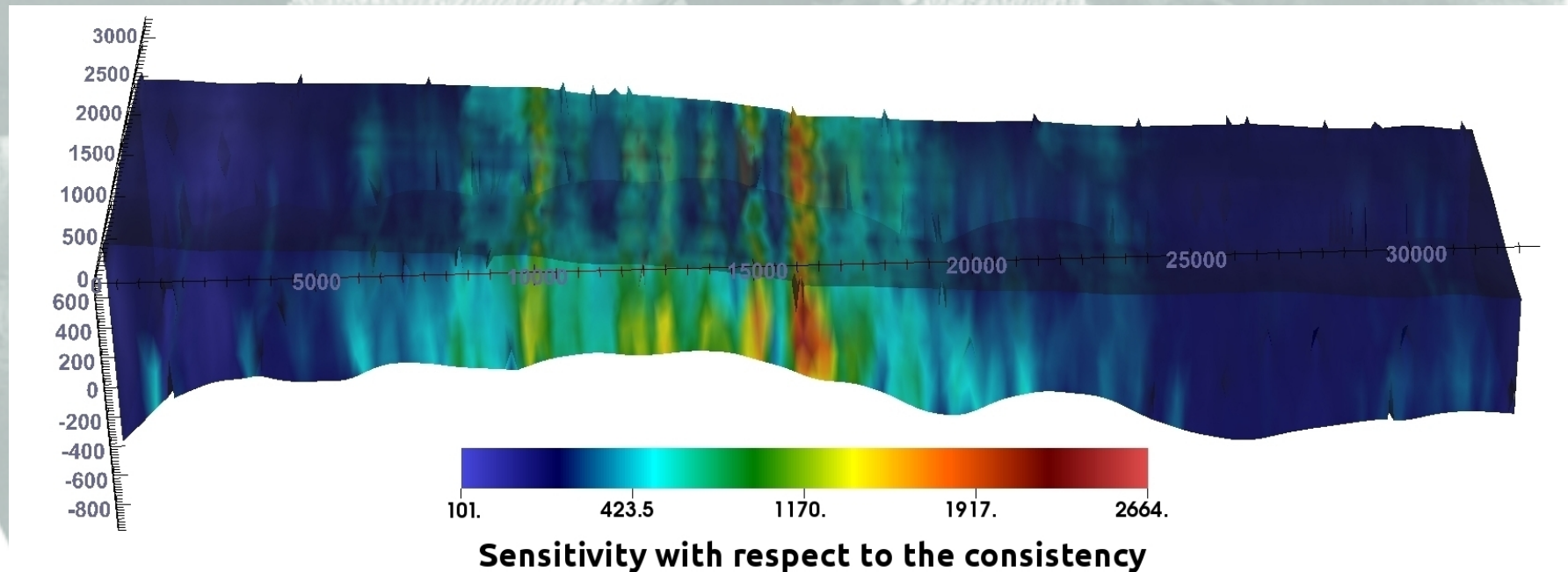
DassFlow-3D operational: tetrahedra, velocity-pressure formulation
(mainly developed by R. Madec – Post-doc IMT, ANR AMAC & ADAGe 2009-2013)
Available online open-source soon



Madec-Martin-Monnier, under preparation

Some perspectives

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Madec-Martin-Monnier, under preparation

Some perspectives

- Introduce the Bingham threshold into LA algorithm
(*a priori* suited for non-differentiable problem)

Glowinski and Wachs, 2011

➔ Herschel-Bulkley constitutive law:
mud flows, avalanches, etc.



Avalanche deposit, Ancy, 2007

- Unsteady adjoint formulation: moving domain (free surface)
shape optimization problem

➔ extra terms due to free-surface movements (*Leibniz formula*)